

Viscous dissipation effect on steady free convective hydromagnetic heat transfer flow of a reactive viscous fluid in a bounded domain

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A study of steady free convective hydromagnetic flow of a reactive viscous fluid in a bounded domain has been analyzed in the present paper. The existence and uniqueness of the solution of momentum and energy equations are proved. By employing the perturbation technique, the momentum and energy equation are solved analytically. Results are presented for velocity and temperature profiles. It is shown that velocity profiles increases as Grashof number, heat source parameter and Prandtl number increases whereas the temperature profile increases as heat source parameter increases.

Keywords: Hydro magnetic flow, Free convection, Steady, Numerical solution

Abstract

1. Introduction

Many researchers carried on MHD (magnetohydrodynamics) flow of the boundary layer problem or a radiating gas inside a vertical pathway. The impact of this type of viscous dissipation term on an unsteady condition was often ignored. The influence of this heat dissipation function cannot be neglected from a practical point of view because of its momentous in several flow issues. It is the bearing lubricant that provides the source of temperature rise and geodynamic heating. For much lower velocity methods the impact of viscous dissipation in the temperature profile is comparatively small. The influence of viscous dissipation cannot be neglected in the manner concerning a dynamic temperature which is analogous to the attributed difference in heat transfer temperature. The boundary layer theory is utilized to analyze the viscous dissipation effect for both incompressible and compressible flows. Magneto hydrodynamics of an electrically conducting fluid is encountered in many problems in geophysics, astrophysics, engineering applications, and other industrial areas. Hydromagnetic free convection flows have a great significance for the applications in the field of stellar and planetary magnetospheres and aeronautics. Engineers employed magneto hydrodynamics principles in the design of heat exchangers pumps, in space vehicle propulsion, thermal protection, control and re-entry, and in creating novel power generating systems. Hydro magnetic flow and heat transfer problems have also become more important, industrially. In many metallurgical processes involving the cooling of many continuous strips of filaments by drawing them through an electrically conducting fluid subject to a magnetic field, the rate of cooling can be controlled and final product of desired characteristics can be achieved. Another important application of Hydromagnetic to metallurgy lies in the purification of molten metals from non-metallic inclusions by the application of a magnetic field [1-16].

In recent decades, more importance have been made to the topic of viscous dissipation effect on unsteady mixed convection, because of its utilitarian aspects in diverse fields as medical chemistry medicine engineering, industries etc. The effects of viscous dissipation on convective heat transfer is significant especially for higher velocity flows, highly viscous flows even at moderate velocities, for fluids with a moderate Prandtl number and moderate velocities with small wall-to-fluid temperature difference or with low wall heat fluxes. Free convection flow of heat and mass transfer problems has diverse applications in industrial processes. The Influence of a magnetic field on free convection flow problems has many applications in science and technology. In the study of MHD flow, the influence of free convection plays a vital role in the fields of petroleum, engineering and agriculture. The final result of the magnetic field and thermal radiation on the induced magnetic field has attracted many researchers. It has been adopted in various fields such as aerodynamics, science and technology, astrophysics and geophysics. In polymer processing industry radiation effects play a major role in getting the final product, by adjusting the heat and mass transfer. The impact of heat and mass transfer plays a major role in distribution of moisture and temperature in agriculture, modeling the chemical equipment, emergence of fog, harm of a crop due to freeze, pollution in the ecosystem, etc. [17-33].

The natural convection flow between heated vertical plates is a classical problem that occurs in many physical phenomena and engineering applications and has received attention in recent years because of its various applications in the design of nuclear reactors, aircraft cabin insulation, thermal storage systems and cooling of electronic equipments. The study of heat transfer and flow of viscous fluids through and across a porous medium has wide ranging applications in various fields of science and engineering. As a result of its technological import to geothermal and reservoir engineering and cooling of nuclear reactors, etc. several researchers have studied such problems in channels, composed of porous materials [34-47].

From the existing literature most of the previous studies considered only the case of free convection with heat transfer, but little has been done in the direction of reactive flow. In this work we shall consider a steady free convective hydro magnetic flow of a reactive viscous fluid in a bounded domain. In section two we provide the mathematical formulation of the problem and method of solution. In section three gives the properties of solution while section four gives the description of the numerical method employed and section five deals with discussion of results.

2. Mathematical formulation and Solution of the problem

Consider a steady one-dimensional free convection flow of an incompressible electrically conducting viscous fluid on a finite plate and the temperatures of the flow assume Arrhenius dependence. The x -axis is along the plate in the upward direction and y -axis normal towards it. A transverse constant magnetic field is assumed to be negligible. By assuming a very small magnetic Reynolds number the induced magnetic field is neglected. The appropriate governing equation is given as:

$$\frac{dv'}{dy'} = 0 \quad (1)$$

$$v' \frac{dv'}{dy'} = \nu \frac{d^2 u'}{dy'^2} + g\beta'(T' - T'_\infty) - \frac{\sigma B_0^2}{\rho} u' - \frac{\nu}{\rho} u' \quad (2)$$

$$v' \frac{dT'}{dy'} = \frac{k}{\rho C_p} \frac{d^2 T'}{dy'^2} - \mu \left(\frac{du'}{dy'} \right)^2 - \frac{Q_0}{\rho C_p} (T' - T'_\infty) \quad (3)$$

The appropriate boundary conditions are:

$$u(y) = 0, T(y) = T_\infty \quad \text{at } y = \pm 1 \quad (4)$$

Introducing the following non-dimensional quantities and assume $v_0 = 0$

$$u = \frac{u'}{u_0}, t = \frac{t' u_0^2}{\nu}, v = \frac{v'}{v_0}, \eta = \frac{u_0 y'}{\nu}, \theta = \frac{T' - T'_\infty}{T'_w - T'_\infty} \quad (5)$$

Equations (2) and (3) one gets the following non-dimensional equations governing the flow:

$$\frac{d^2 u}{d\eta^2} - (M^2 + K)u = -Gr\theta \quad (6)$$

$$\frac{d^2 \theta}{d\eta^2} + \text{Pr} E \left(\frac{du}{d\eta} \right)^2 - Q \text{Pr} \theta = 0 \quad (7)$$

where

$$E = \frac{\rho u_0^2}{(T'_w - T'_\infty)}, Gr = \frac{\rho \beta' (T'_w - T'_\infty) \nu}{u_0^3}, Q = \frac{Q_0}{\rho C_p u_0^2}, \text{Pr} = \frac{\mu C_p}{k}, M^2 = \frac{B_0^2 \nu \sigma}{\rho u_0^2}$$

The corresponding boundary condition in dimensionless form are reduced to

$$u = 0, \theta = 0 \quad \text{at } \eta = \pm 1 \quad (8)$$

The physical variables u, θ and C can be expanded in the power of Eckert number (E). This can be possible physically as E for the flow of an incompressible fluid is always less than unity. It can be interpreted physically as the flow due to the Joules dissipation is super imposed on the main flow. Hence we can assume

$$\begin{aligned} u(\eta) &= u_0(\eta) + Eu_1(\eta) + O(E^2) \\ \theta(\eta) &= \theta_0(\eta) + E\theta_1(\eta) + O(E^2) \end{aligned} \quad (9)$$

Using equation (9) in equations (6)–(7) and equating the coefficient of like powers of E , we have

$$u_0'' - M^2 u_0 = Gr\theta_0 \quad (10)$$

$$u_1'' - M^2 u_1 = Gr\theta_1 \quad (11)$$

$$\theta_0'' - QPr\theta_0 = 0 \quad (12)$$

$$\theta_1'' - QPr\theta_1 = -Pr u_0'^2 \quad (13)$$

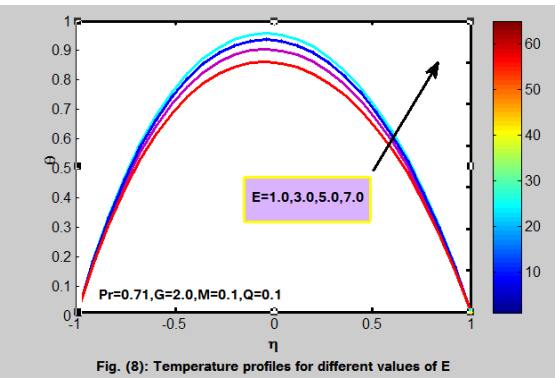
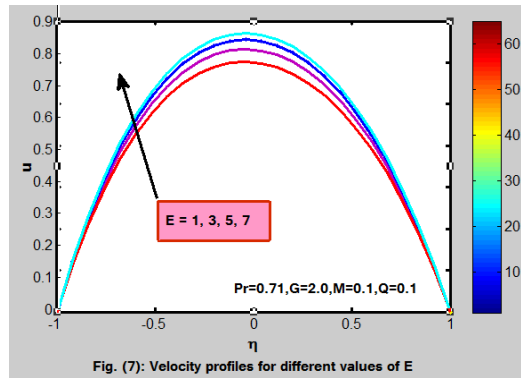
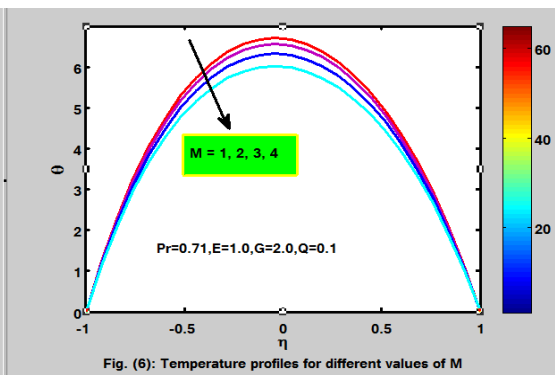
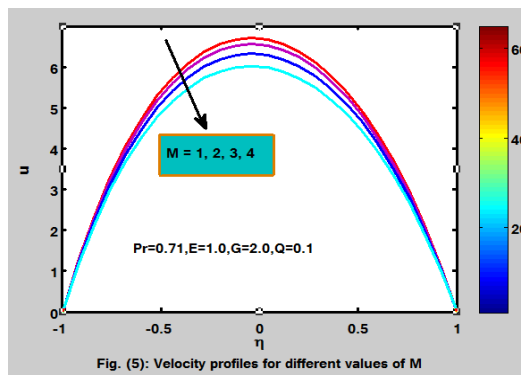
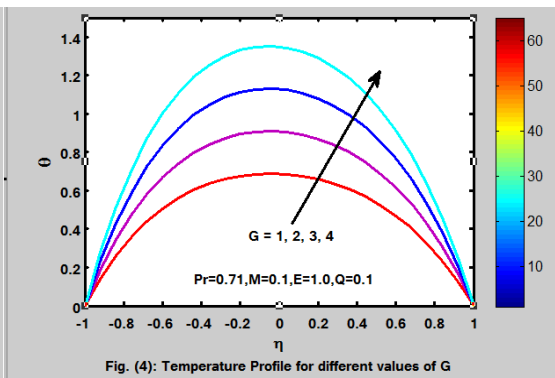
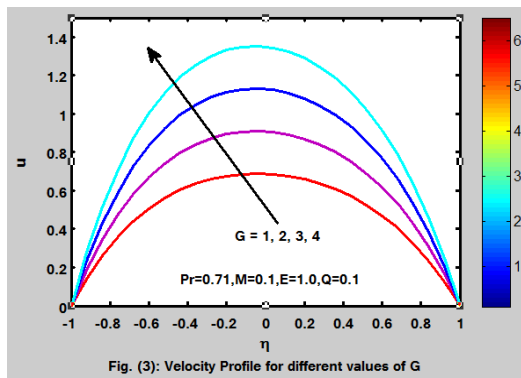
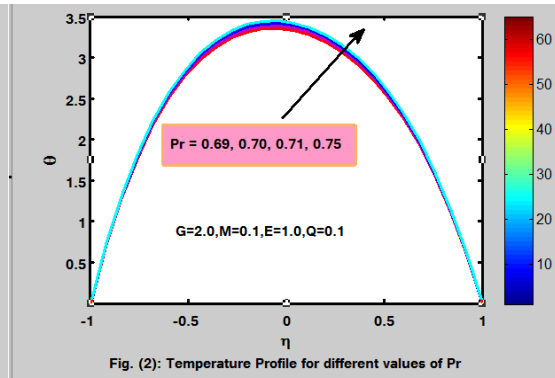
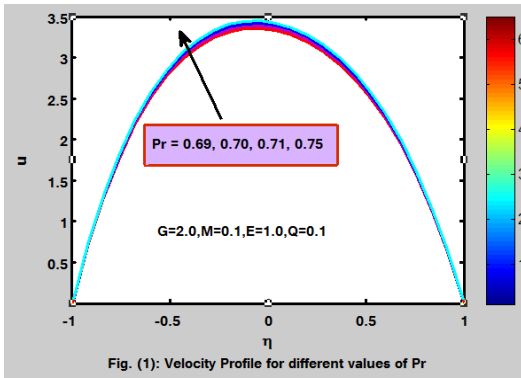
The corresponding boundary condition in dimensionless form are reduced to

$$u_0 = 0, \theta_0 = 0, u_1 = 0, \theta_1 = 0 \quad \text{at } \eta = \pm 1 \quad (14)$$

Using equations (10) – (12) under the boundary conditions (14), we obtain the solution of the velocity and temperature.

3. Results and Discussion

In this section, numerical computations are presented in the form of non-dimensional velocity and temperature profiles. Also, numerical computation has been carried out for different values of the parameters entering into the problem. The values of Grashof number are taken to be large from the physical point of view. The large Grashof number corresponds to the cooling problem. Figures (1) and (2) reveal the effect of Prandtl number on velocity and temperature profiles. High thermal conductivity and are therefore a good choice for heat conducting liquids. As can be seen, liquid metals are very good heat transfer liquids. Interestingly, air is a decent heat transfer liquid as, well, whereas typical organic solvents are not. It is observed that the velocity increases while temperature decreases as the Prandtl number increases. Figures (3) and (4) showed that the velocity increases and the temperature do not change as Grashof number increases. Figures (5) and (6) showed the effect of Hartmann number on the velocity as well as temperature profiles. It is observed that the velocity decrease as the Hartmann number increases; while Hartmann number does not have noticeable effect on the temperature profile. Figures (7) and (8) reveal that Grashof number and activation energy parameter does not have noticeable effect on both velocity and temperature profiles respectively. It is shown that both velocity and temperature profiles increases as Frank-Kamenetskii parameter increases. Finally, it is observed in all that the flow is symmetric about the centre (i.e., $\eta = 0$).



Conclusion

In this paper we derived the energy and momentum equations governing free convective hydro magnetic flow of a reactive viscous fluid. We have shown the existence and uniqueness of solutions, and then presented numerical results for various non–dimensionless parameters graphically. From the present study we can make the following conclusions:

- The velocity profiles increases whereas temperature profile does not have noticeable effect with an increase of the free convection current.
- Using magnetic field we can control the flow characteristic and heat transfer.
- The velocity and temperature profiles increase as the Frank-Kamenetskii parameter increases.
- The velocity profile increases whereas the temperature profile decreases as the Prandtl number increases.

Nomenclature:	
Gr	Grashof number
Pr	Prandtl number
M^2	Magnetic parameter
K	Porous permeability
E	Eckert number
Q	Heat source parameter

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